

Non-ideal effects on the typical trailing edge shock pattern of ORC turbine blades

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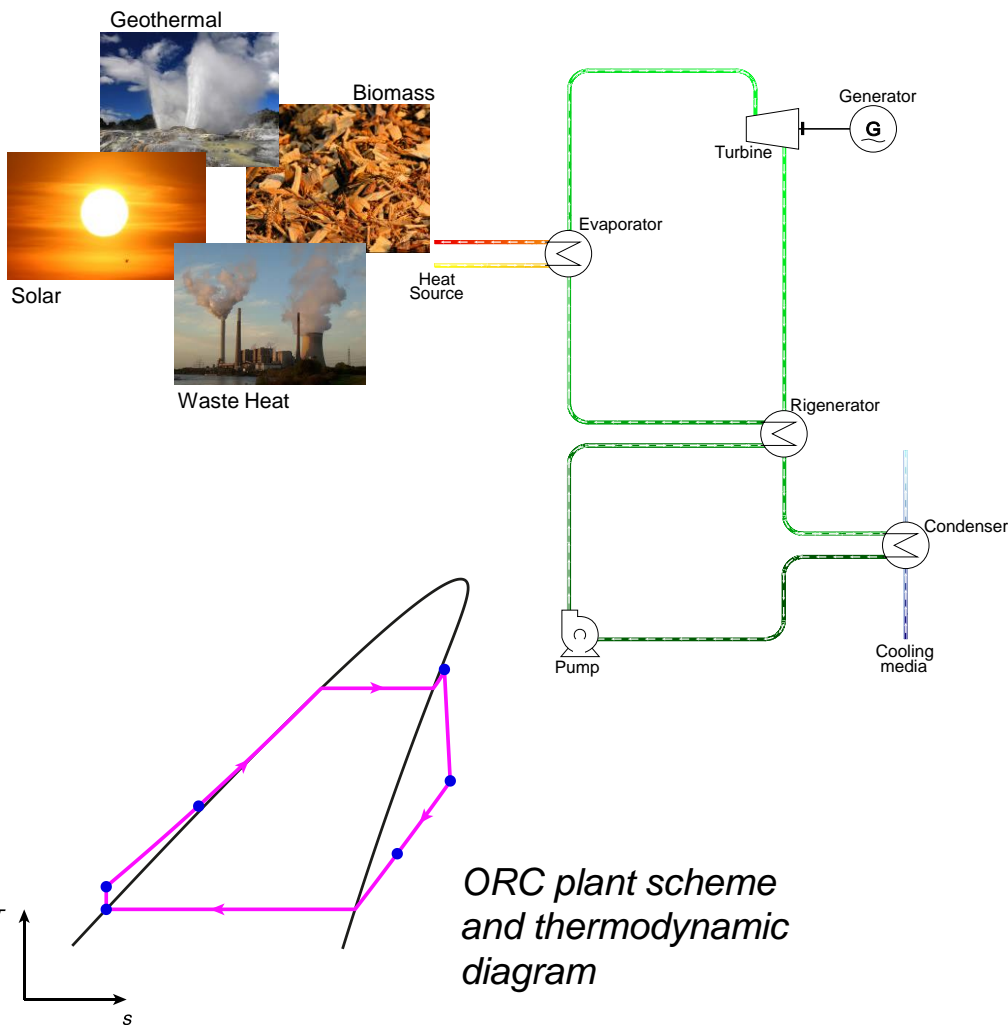
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MOTIVATIONS



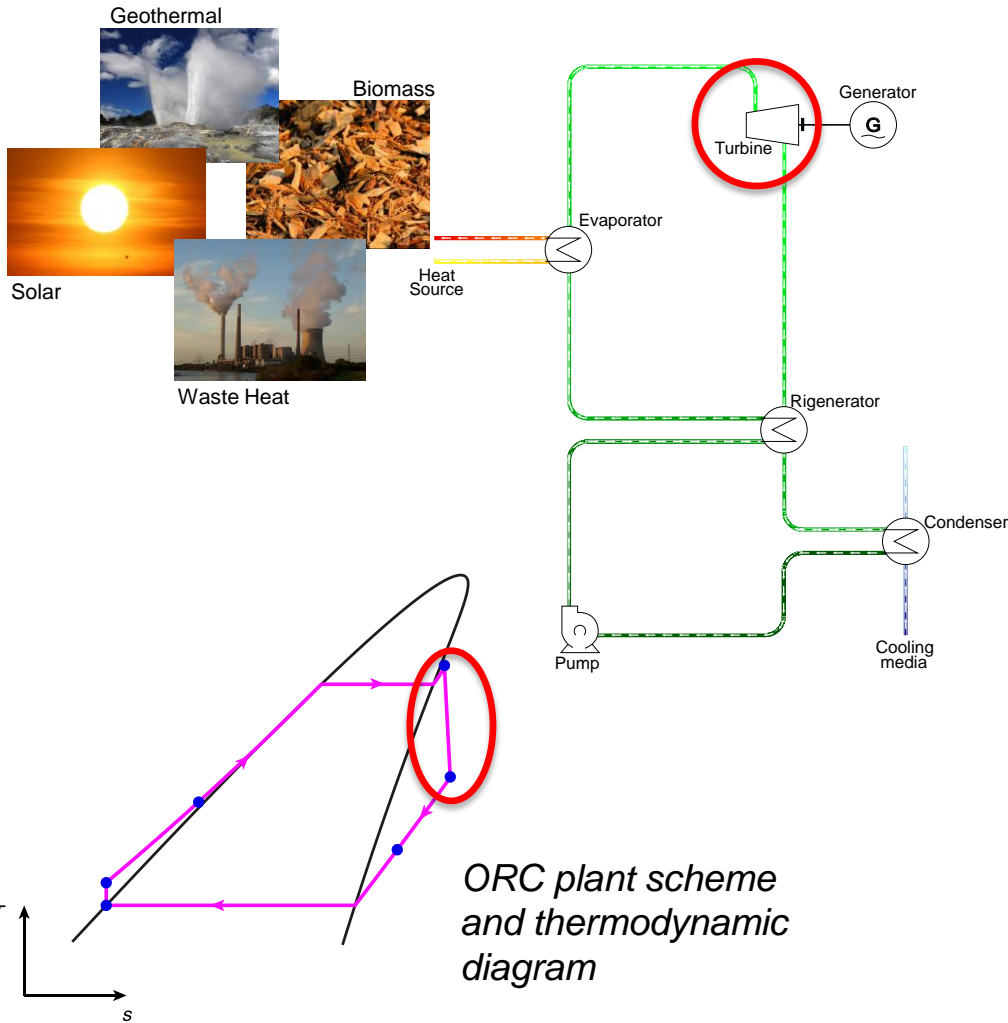
ORC attractive features

- Adaptability to various (low temperature) heat sources
- Lower complexity wrt steam cycle
- Turbine technical advantages wrt steam turbine (lower rpm, lower pressures, no erosion)
- High flexibility
- ...

ORC challenges

- Choice of suitable working fluid
- Transient phenomena
- Complex thermodynamic modelling of the working fluid
- Heat exchangers and turbine design
- ...

MOTIVATIONS



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ORC challenges

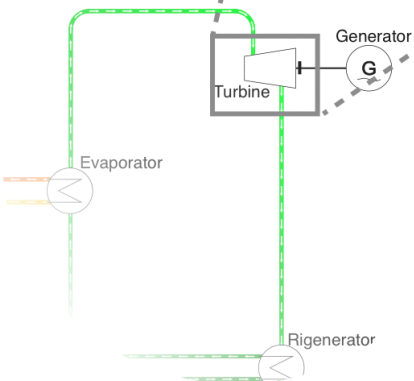
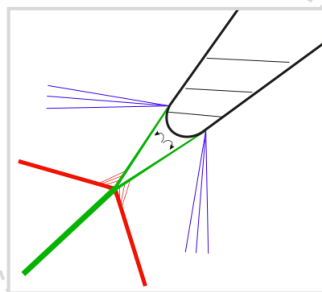
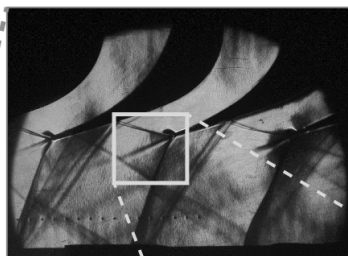
- Choice of suitable working fluid
- Transient phenomena
- **Complex thermodynamic modelling of the working fluid**
- Heat exchangers and **turbine design**
- ...

MOTIVATIONS

ORC turbine

- Typically few stages (often one only)
- High pressure ratio
- Design expansion through the non-ideal regime: low values of the speed of sound → highly supersonic flow
- Shock waves: fish-tail shocks, post-expansion, off-design
- Large contribution of inviscid loss to total loss

Flow in transonic ORC vanes



Research question

How do non-ideal effects across oblique shocks impact on the design of ORC turbines?

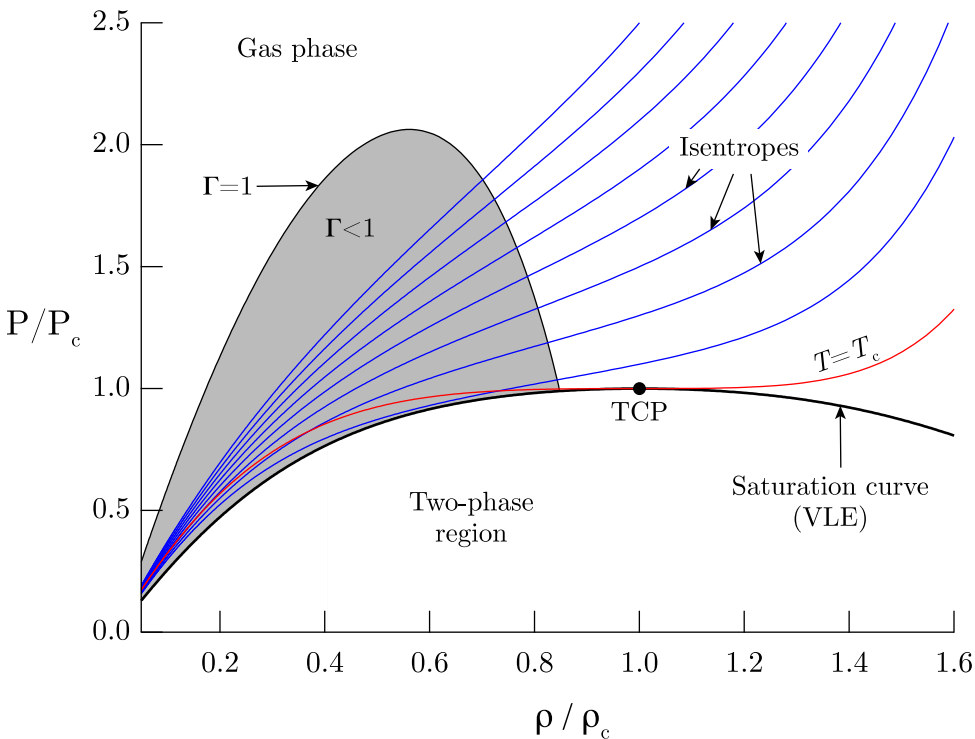
PRESENTATION OUTLINE

- Introduction: NICFD
- Methodology
- Oblique shocks in the non-ideal flow regime
- Application: oblique shocks in siloxane MDM
- Discussion and concluding remarks



Non-Ideal Compressible Fluid Dynamics

$$Pv \neq RT$$



Fluid: MDM (RefProp)

Features

- Subject: dense vapours, supercritical fluids, two-phase compressible flows
- Compressibility
- Phase transition
- Critical point

Application

- ORC
- Supercritical CO₂
- Refrigeration
- Oil & Gas compression/expansion
- ...

Non-Ideal Compressible Fluid Dynamics

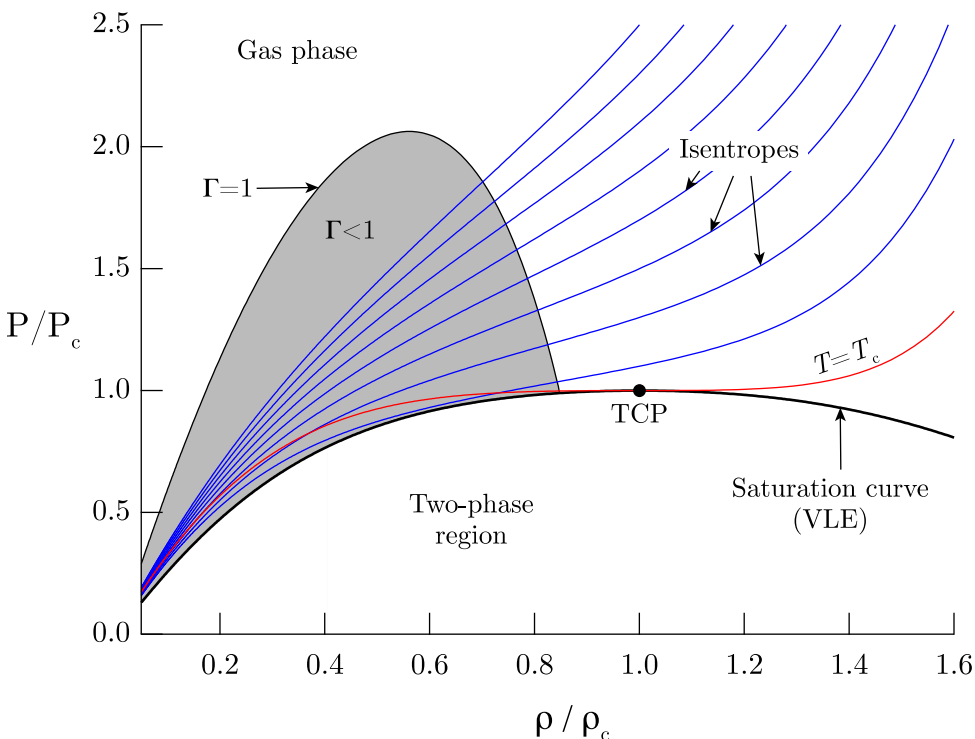
$$\parallel$$

$$Pv \neq RT$$

Measure of non-ideality in compressible flows:

The fundamental derivative of gasdynamics
(Thompson 1971)

$$\Gamma = 1 + \frac{\rho}{c} \left(\frac{\partial c}{\partial \rho} \right)_s = 1 + \frac{\rho}{c} \left(\frac{\partial^2 P}{\partial \rho^2} \right)_s = \frac{v^3}{2c^2} \left(\frac{\partial^2 P}{\partial v^2} \right)_s$$



Fluid: MDM (RefProp)

Gasdynamic regimes

- $\Gamma > 1$ Ideal regime
- $\Gamma < 1$ Non-Ideal regime
 - $0 < \Gamma < 1$ Non-Ideal classical regime
 - $\Gamma < 0$ Non-Classical regime

Rankine-Hugoniot relations

$$h_A - \frac{1}{2}P_A(v_A + v_B) = h_B - \frac{1}{2}P_B(v_A + v_B)$$

$$\sqrt{-\frac{(P_B - P_A)}{(v_B - v_A)}} = \rho_A |\mathbf{u}_A| \sin \beta$$

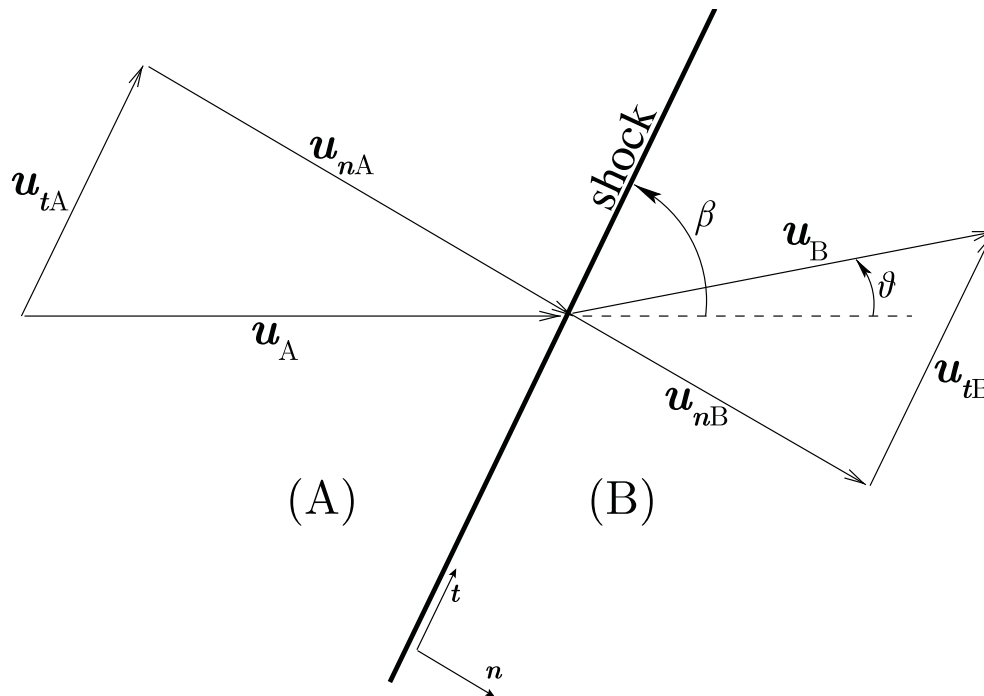
$$\rho_A \tan \beta = \rho_B \tan(\beta - \theta)$$

$$|\mathbf{u}_A| \cos \beta = |\mathbf{u}_B| \cos(\beta - \theta)$$

Admissibility conditions

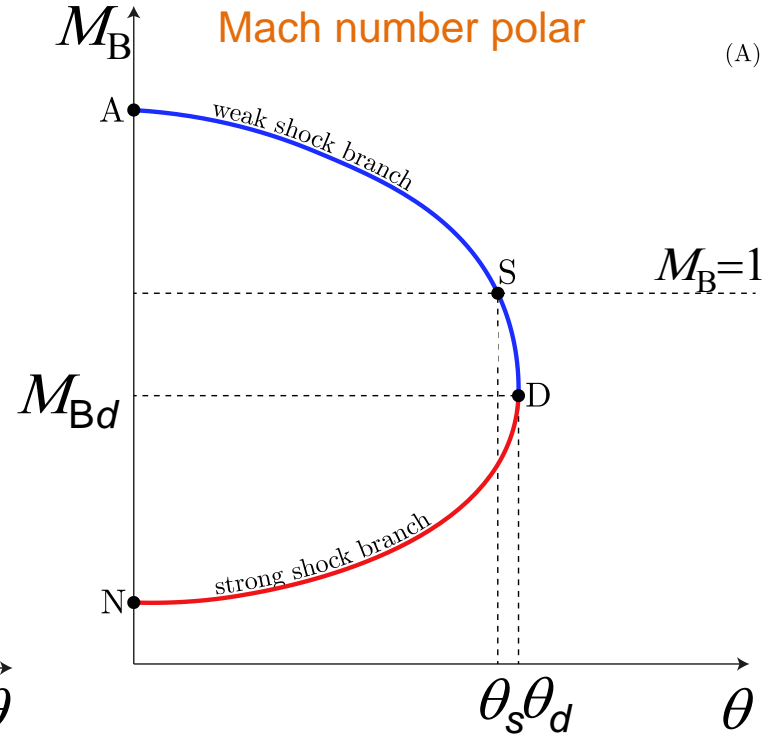
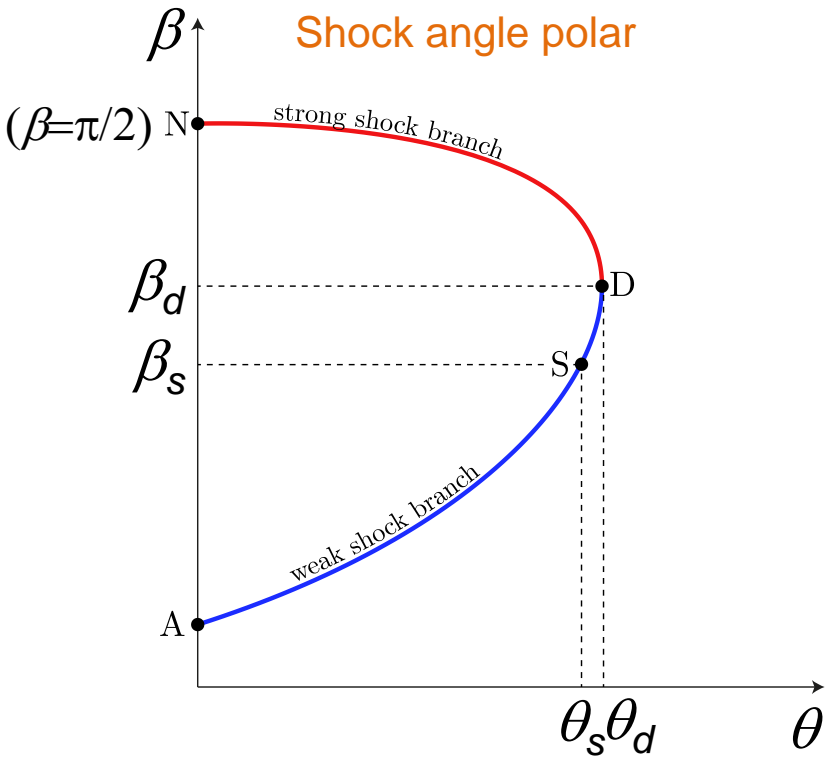
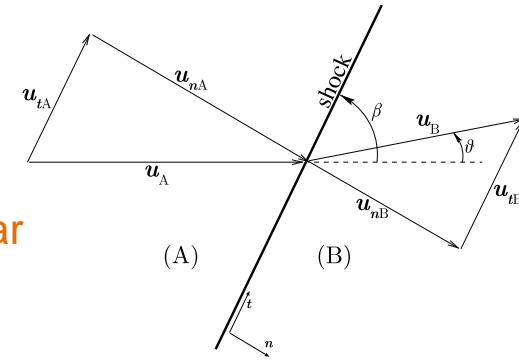
$$s_B > s_A$$

$$\left(M_{nB} = \frac{|\mathbf{u}_{nB}|}{c_B} < 1 < \frac{|\mathbf{u}_{nA}|}{c_A} = M_{nA} \right)$$



METHODOLOGY

Deflection shock polars: X- θ diagrams for fixed upstream state



A – Mach wave (acoustic limit)

N – normal shock

D – detachment point (max deflection)

S – downstream sonic point

OBLIQUE SHOCKS IN THE NON-IDEAL REGIME

Perfect-gas: explicit formulas

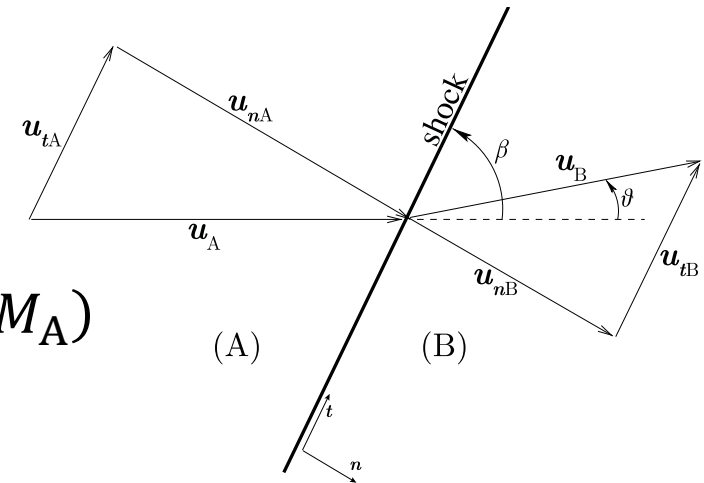
$$\tan \theta = \frac{2}{\tan \beta} \left[\frac{M_A^2 \sin^2 \beta - 1}{M_A^2 (\gamma + \cos 2\beta) + 2} \right]$$

$\beta(\vartheta, M_A)$

$$\frac{\rho_B}{\rho_A} = \frac{(\gamma + 1) M_A^2 \sin^2 \beta}{2 + (\gamma - 1) M_A^2 \sin^2 \beta}$$

$$\frac{P_B}{P_A} = 1 + \frac{2\gamma}{\gamma + 1} (M_A^2 \sin^2 \beta - 1)$$

$$M_B^2 = \frac{1}{\sin^2(\beta - \vartheta)} \frac{1 + \frac{\gamma - 1}{2} M_A^2 \sin^2 \beta}{\gamma M_A^2 \sin^2 \beta - \frac{\gamma - 1}{2}}$$



Dependencies:

- Deflection angle ϑ
- Upstream Mach number M_A

No dependence on the upstream thermodynamic state (e.g. P_A, ρ_A)

OBLIQUE SHOCKS IN THE NON-IDEAL REGIME

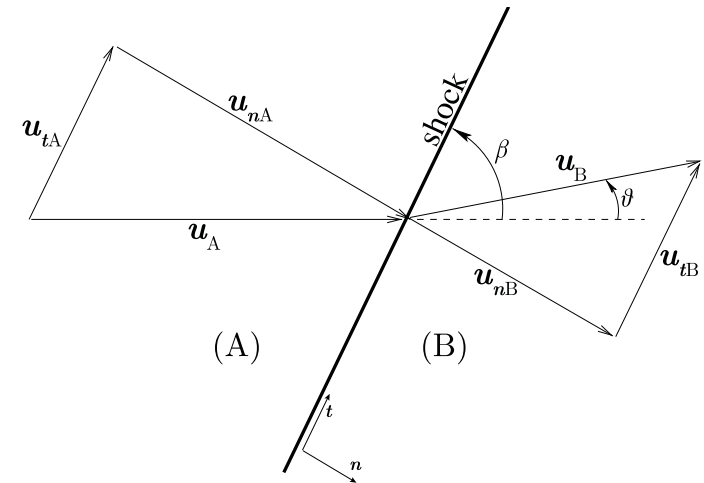
Non-ideal regime: acoustic limit

$$\beta = \sin^{-1}(1/M_A) + \frac{\Gamma_A}{2} \frac{M_A^2}{M_A^2 - 1} \vartheta + \mathcal{O}(\vartheta^2)$$

$$\frac{\rho_B}{\rho_A} = 1 + \frac{M_A \Gamma_A}{\sqrt{M_A^2 - 1}} \vartheta + \mathcal{O}(\vartheta^2)$$

$$\frac{P_B}{P_A} = 1 + \frac{\rho_A c_A^2}{P_A} \frac{M_A \Gamma_A}{\sqrt{M_A^2 - 1}} \vartheta + \mathcal{O}(\vartheta^2)$$

$$M_B = M_A + \left(1 - \Gamma_A - \frac{1}{M_A^2}\right) \frac{M_A^3}{\sqrt{M_A^2 - 1}} \vartheta + \mathcal{O}(\vartheta^2)$$



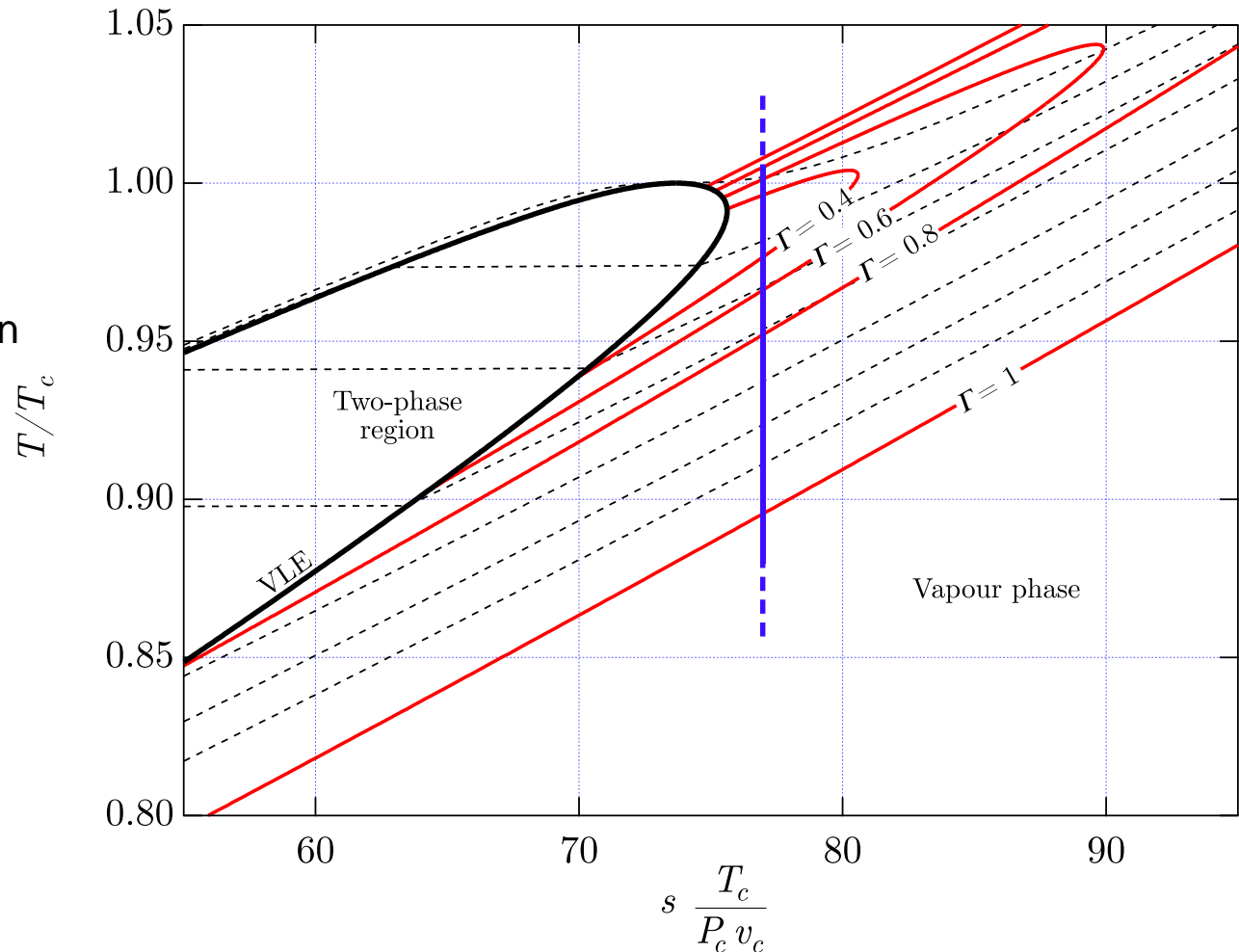
Dependences:

- Deflection angle ϑ
- Upstream Mach number M_A
- Upstream thermodynamic state (e.g., P_A , ρ_A)

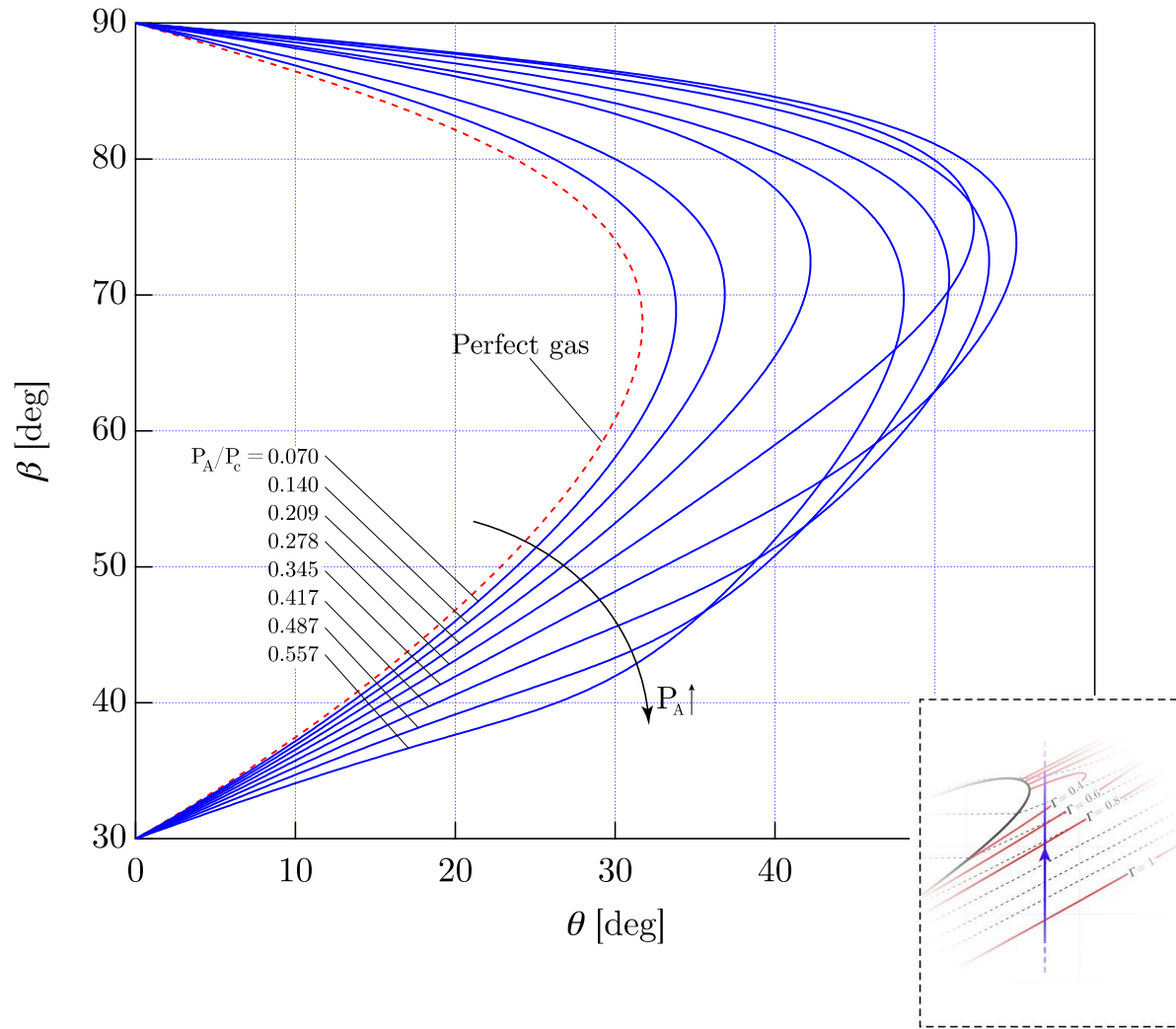
APPLICATION: OBLIQUE SHOCKS IN MDM

Parametric study

- Fluid: siloxane MDM (RefProp)
- Fixed upstream entropy in the non-ideal thermodynamic regime ($\Gamma < 1$)
- Fixed upstream Mach number ($M_A = 2$)



APPLICATION: OBLIQUE SHOCKS IN MDM



$\beta - \vartheta$ Diagram

- Acoustic limit:

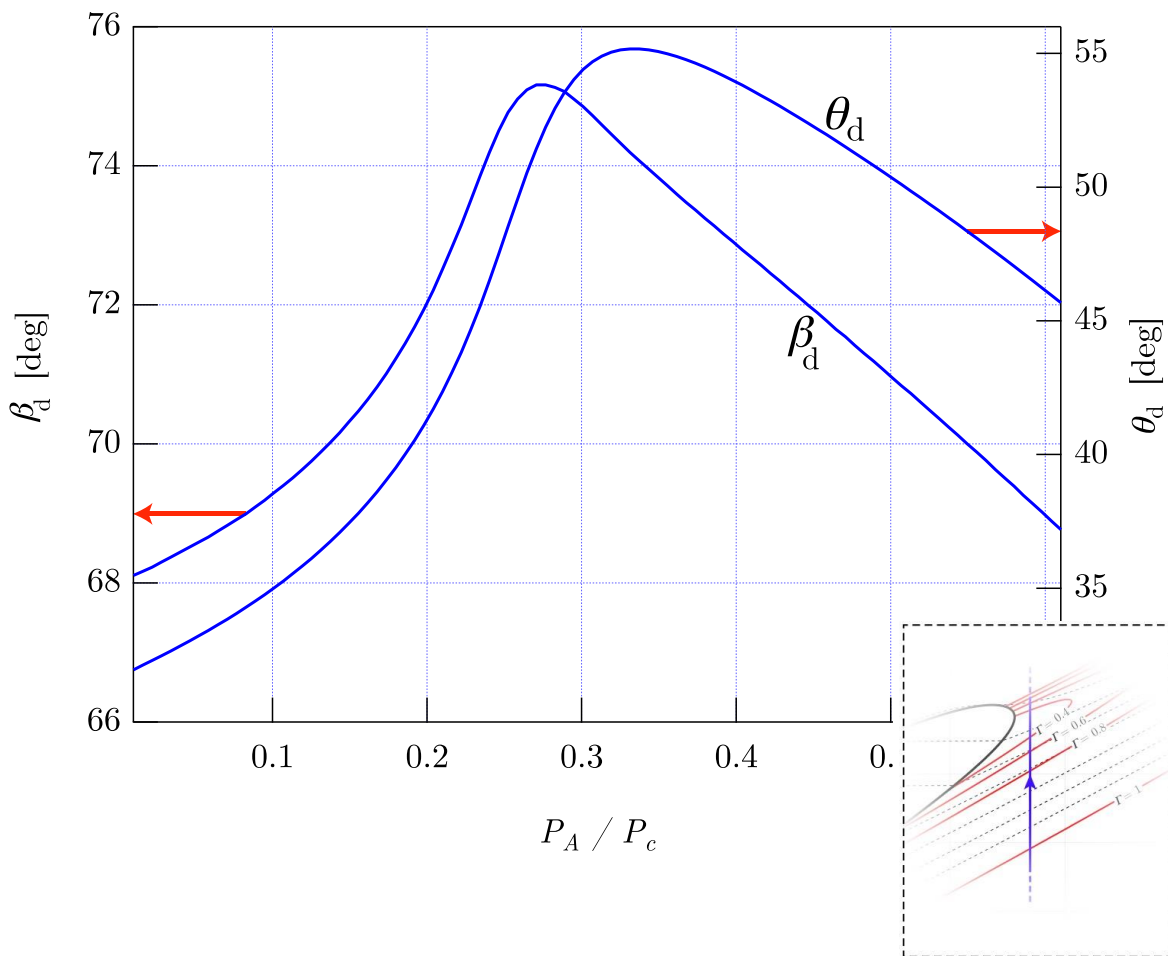
$$\beta = \sin^{-1}(1/M_A) + \frac{\Gamma_A}{2} \frac{M_A^2}{M_A^2 - 1} \vartheta$$

dependence on the upstream tmd state through Γ_A

- Strong dependence of the detachment angles on the upstream tmd state

APPLICATION: OBLIQUE SHOCKS IN MDM

$\beta - \vartheta$ Diagram



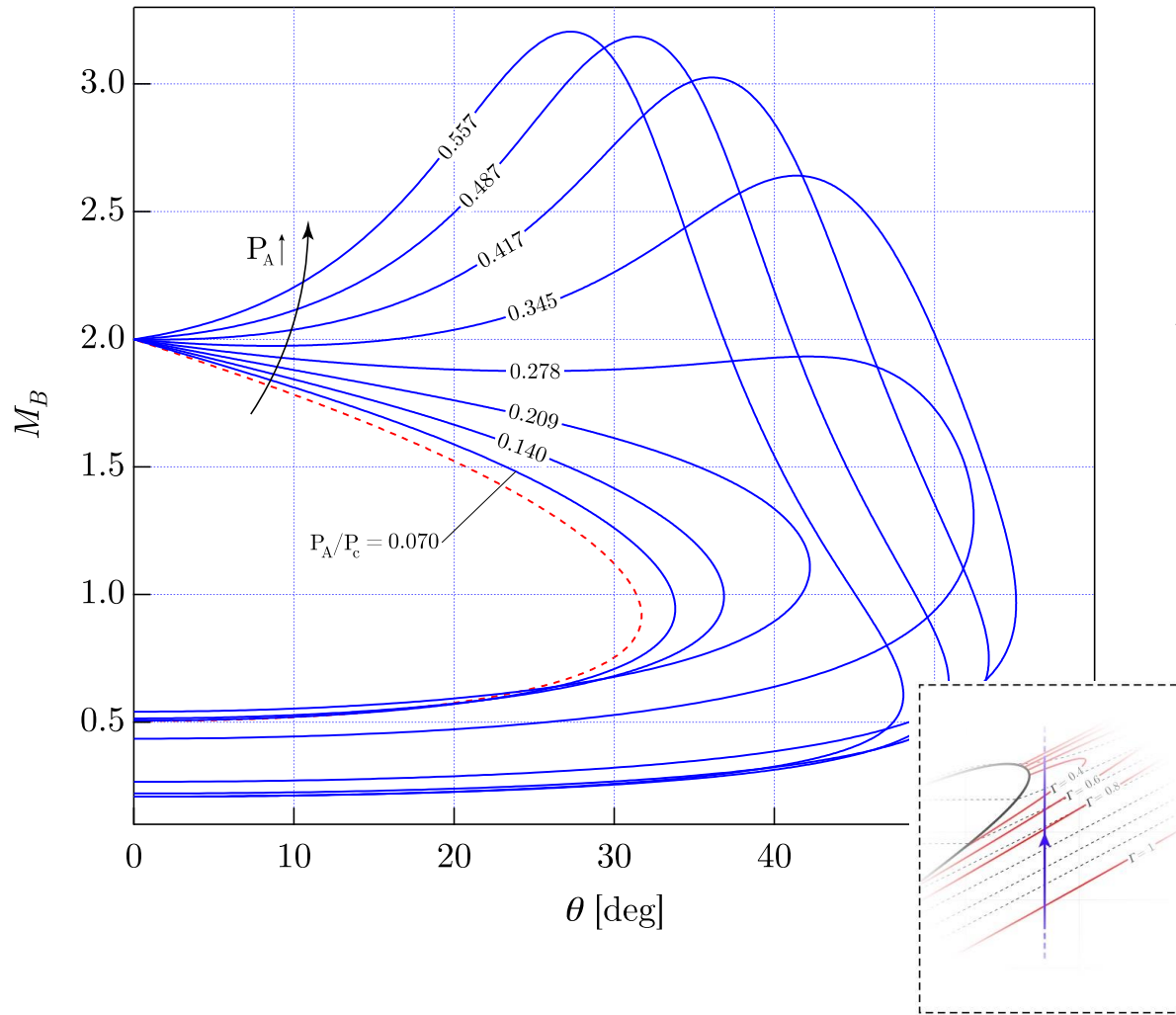
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dependence on the upstream
tmd state through Γ_A

- Strong dependence of the
detachment angles on the
upstream tmd state

APPLICATION: OBLIQUE SHOCKS IN MDM



$M_B - \vartheta$ Diagram

- Acoustic limit:

$$M_B = M_A + \left(1 - \Gamma_A - \frac{1}{M_A^2}\right) \frac{M_A^3}{\sqrt{M_A^2 - 1}} \vartheta$$

dependence on the upstream tmd state through Γ_A

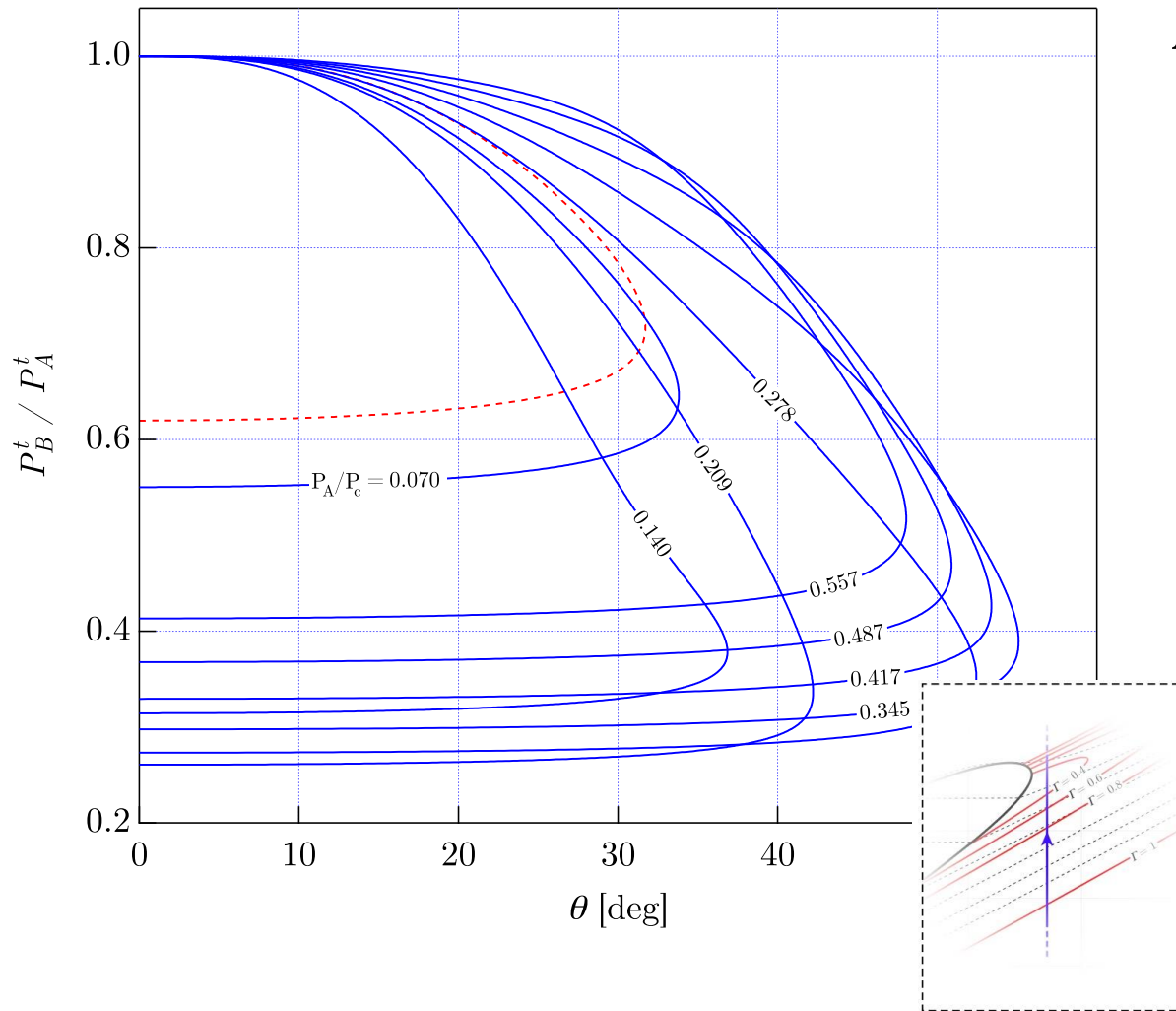
- Non-ideal oblique shocks (Mach number increasing):

$$\Gamma_A < 1 - \frac{1}{M_A^2}$$

for small deviations

APPLICATION: OBLIQUE SHOCKS IN MDM

$P_B^t / P_A^t - \vartheta$ Diagram



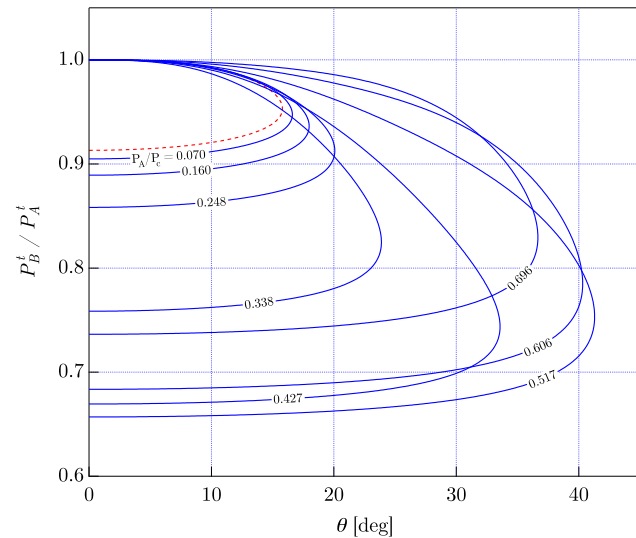
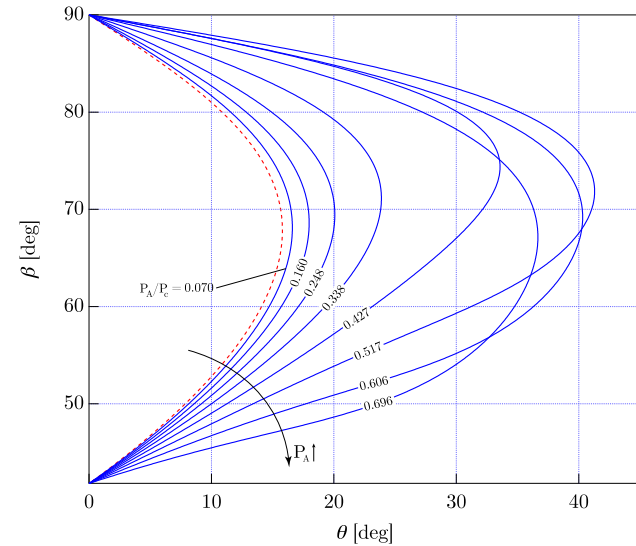
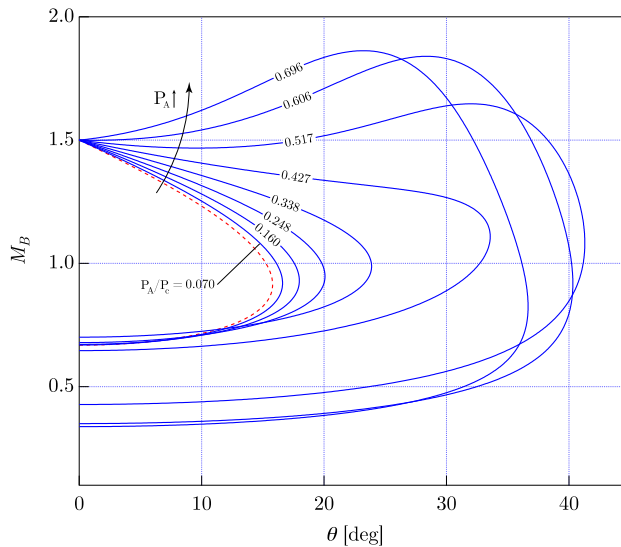
- Fixed $\vartheta \rightarrow$ non-monotonic variation of the shock loss with the upstream pressure
- Larger shock loss across strong oblique shocks w.r.t perfect-gas case
- Smaller shock loss across weak oblique shocks w.r.t perfect-gas case for low values of Γ_A

APPLICATION: OBLIQUE SHOCKS IN MDM

Further parametric study

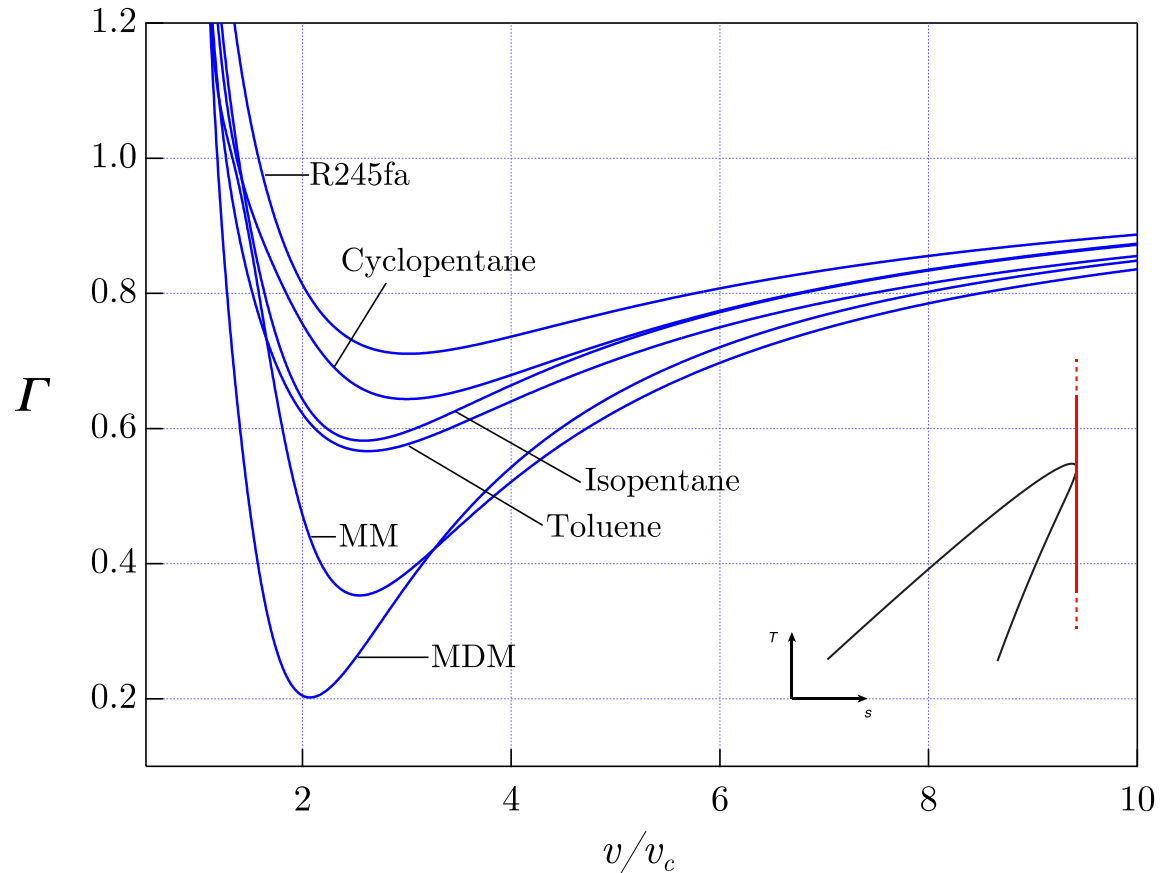
- Same fluid, same isentrope, $M_A = 1.5$
- Non-ideal oblique shocks (Mach number increasing): lower threshold on the upstream Mach number

$$M_{A,\min} = 1/\sqrt{1 - \Gamma_{A,\min}}$$



Extension to other fluids

- Same qualitative behaviour expected for most moderate-to-high molecularly complex fluids
- Qualitatively similar thermodynamic topology of the fundamental derivative of gasdynamics



Extension to other fluids

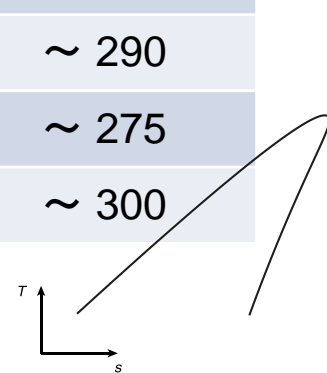
- Non-ideal oblique shocks (Mach number increasing): lower threshold on the upstream Mach number

$$M_{A,\min} = 1/\sqrt{1 - \Gamma_{A,\min}}$$

- Total conditions for non-ideal oblique shocks may exceed thermal stability limit

Example: minimum P^t and T^t for non-ideal oblique shocks along isentrope tangent to VLE

Fluid	P_{\min}^t [bar]	T_{\min}^t [°C]	T_{\lim} [°C]
MDM	16.55	299.0	~ 290÷300
MM	25.90	262.0	~ 300
Toluene	74.00	355.0	~ 400
Isopentane	64.80	221.0	~ 290
Cyclopentane	97.30	280.7	~ 275
R245fa	107.77	204.0	~ 300



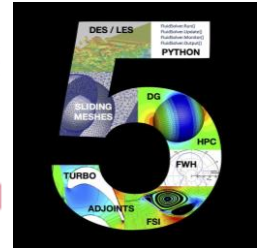
CONCLUSIONS

- Oblique shock waves were investigated in the non-ideal gime because of their relevance to ORC turbine flows
- Main results:
 - Shock angle polar shifts to higher deflection angle
 - Appearance of Mach number-increasing oblique shocks (non-ideal oblique shocks)
 - Shock loss: larger across strong oblique shocks, possibly smaller across weak oblique shocks w.r.t. perfect-gas case
- MDM used for explanatory purposes, direct extension to other molecularly complex fluids employed in ORCs
- Highly non-ideal effects at design conditions only for supercritical ORCs

FUTURE WORK

- Numerical investigation on real vanes configurations at design and off-design conditions
- Experimental observation of non-ideal effects across oblique shock waves at TROVA (Test Rig for Organic Vapours), CREALab PoliMi

SU2
The Open-Source CFD Code



...QUESTIONS?

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